

Gamma-Ray Burst Central Engines *with Binary Population Synthesis Models*

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Gamma-ray burst requirements

Long-duration gamma-ray bursts (GRBs) arise from the collapse of massive, stripped envelope, rapidly spinning stars. We know this because...



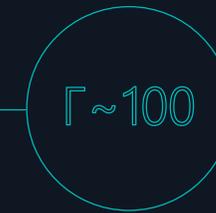
Stripped envelope supernova

They are associated with broad-line type Ic supernovae ($\sim 10^{52}$ erg and lacking in hydrogen and helium) – *observational*



Low metallicity environments

They prefer low-metallicity¹ host environments: implies faster spinning stars – *observational*



Relativistic jets

They launch relativistic jets ($\Gamma \sim 100$), requiring a rotating central engine (see e.g. Blandford-Znajek mechanism) – *theoretical*

What are the central engines?

What powers the launch of relativistic jets in GRBs?

Collapsar

Black hole accretion disc

Narayan et al. 1992,
Woosley et al. 1993

Accretion powered, requires specific angular momentum $> 10^{16} \text{ cm}^2 \text{ s}^{-1}$ at ISCO

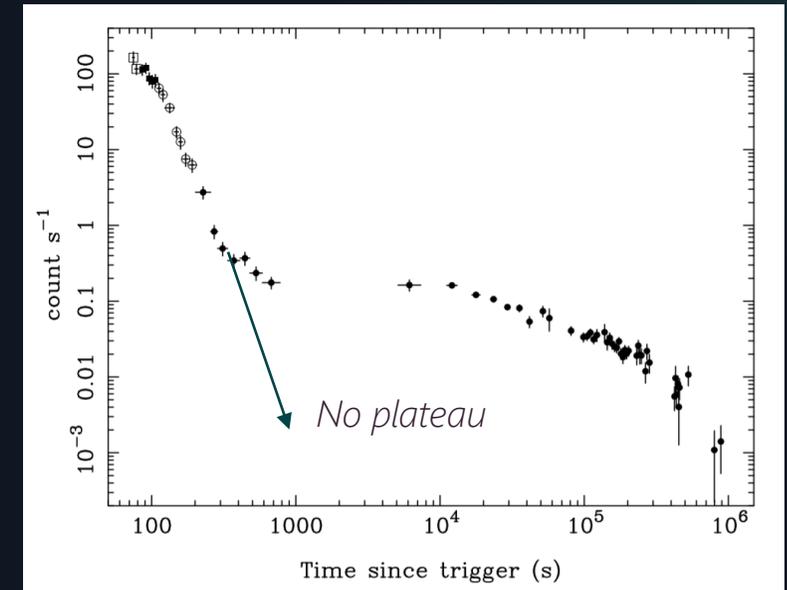
Magnetar

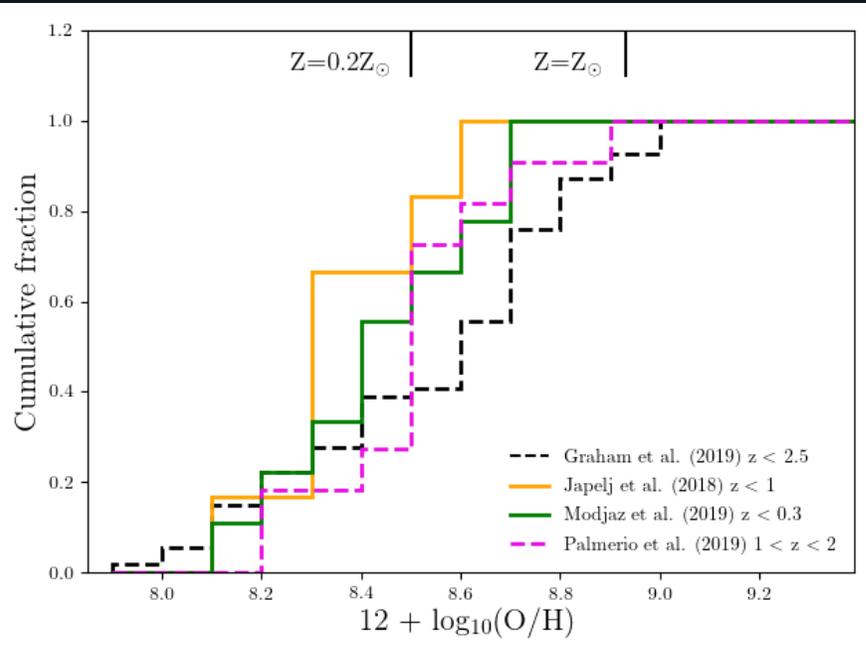
Neutron star magnetic field

Metzger et al. 2011,
Mazalli et al. 2014

E.g. magnetic tower mechanism, requires a ms-period magnetar

Plateaus in GRB X-ray light curves (as in this example from Vaughan et al. 2006) are indicative of continued central engine activity





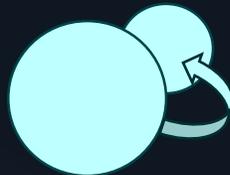
The host galaxy metallicity distribution, as determined by various studies

What are the progenitor systems?

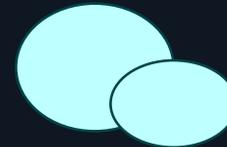
We need a model for long-duration, core-collapse GRBs that explains (i) the GRB prompt and afterglow emission (i.e., works with the collapsar or magnetar models) and (ii) the host galaxy properties.



*Single stars: e.g. Wolf-Rayet, O stars (hard to maintain spin **and** lose envelope)*



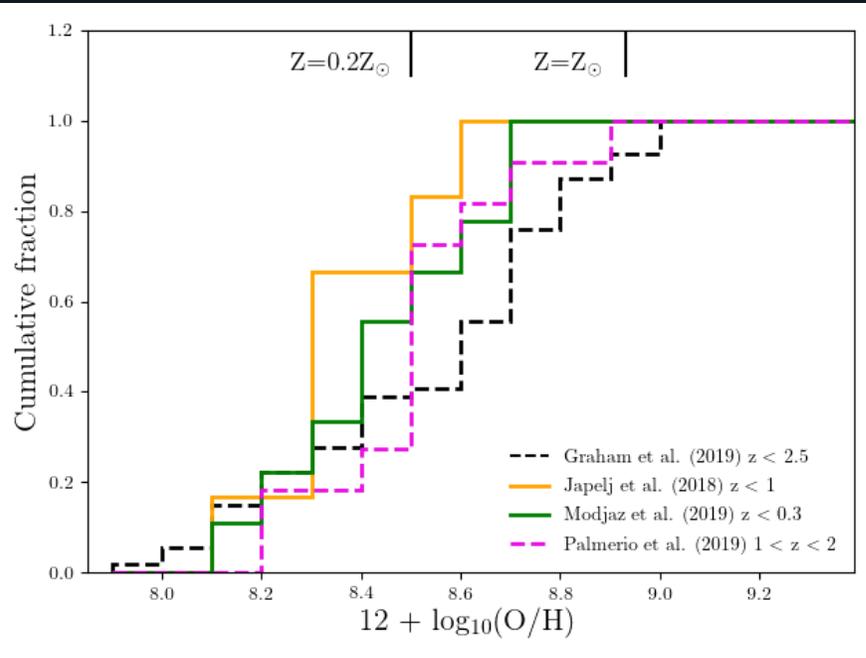
Quasi Homogenous Evolution (QHE): accretion onto a secondary star spins it up to rotational speeds that mix the stellar interior.



Binary stars with tides: tidal interactions maintain spins, counteracting the effects of mass-loss



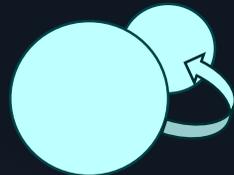
Other: e.g. binary driven hypernova where the GRB is launched by NS \rightarrow BH collapse induced by the SN of a companion



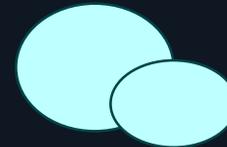
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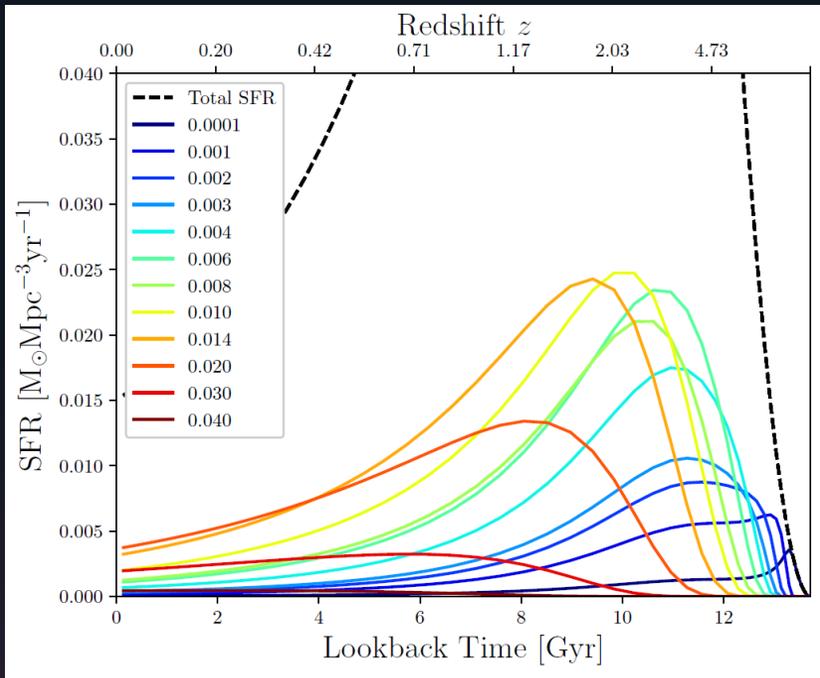
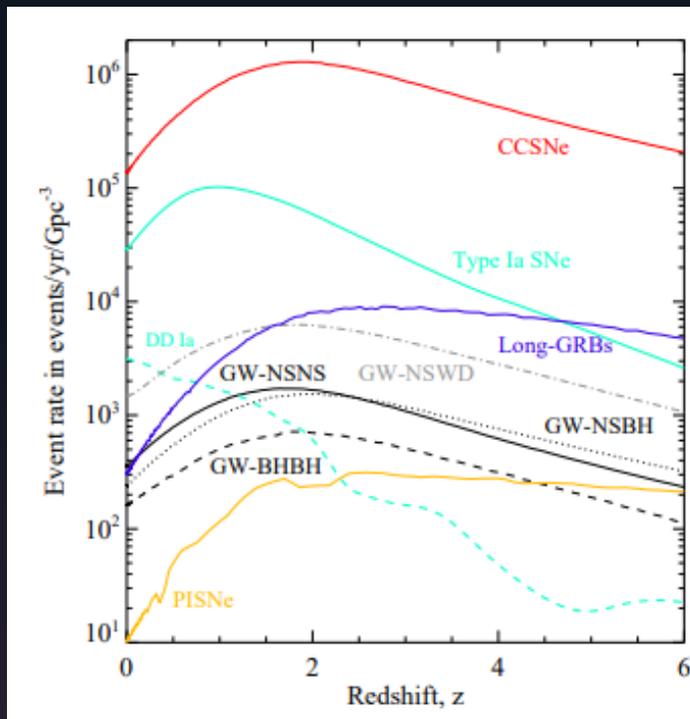


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BPASS

Binary Population and Spectral Synthesis

- Range of models, IMF, binary parameters – see Eldridge et al. 2017 for a full description
- Each model is weighted according to its frequency in a $10^6 M_{\odot}$ stellar population (at the chosen metallicity)
- Project to investigate the plausibility of non-QHE progenitor channels using BPASS – described in Chrimes, Stanway & Eldridge 2019
- Previous studies include Fryer et al. 2007, Trenti et al. 2015...



Using Langer & Norman 2006 distributions.

Eldridge et al. 2019 synthesised the rates of GW and EM transients using BPASS. Only the QHE pathway was considered for long-GRBs.

Total H mass > $5 \times 10^{-4} M_{\odot} \rightarrow$ Type II

If type I:
 He ejecta fraction > 0.2 \rightarrow Type Ib

Model selection

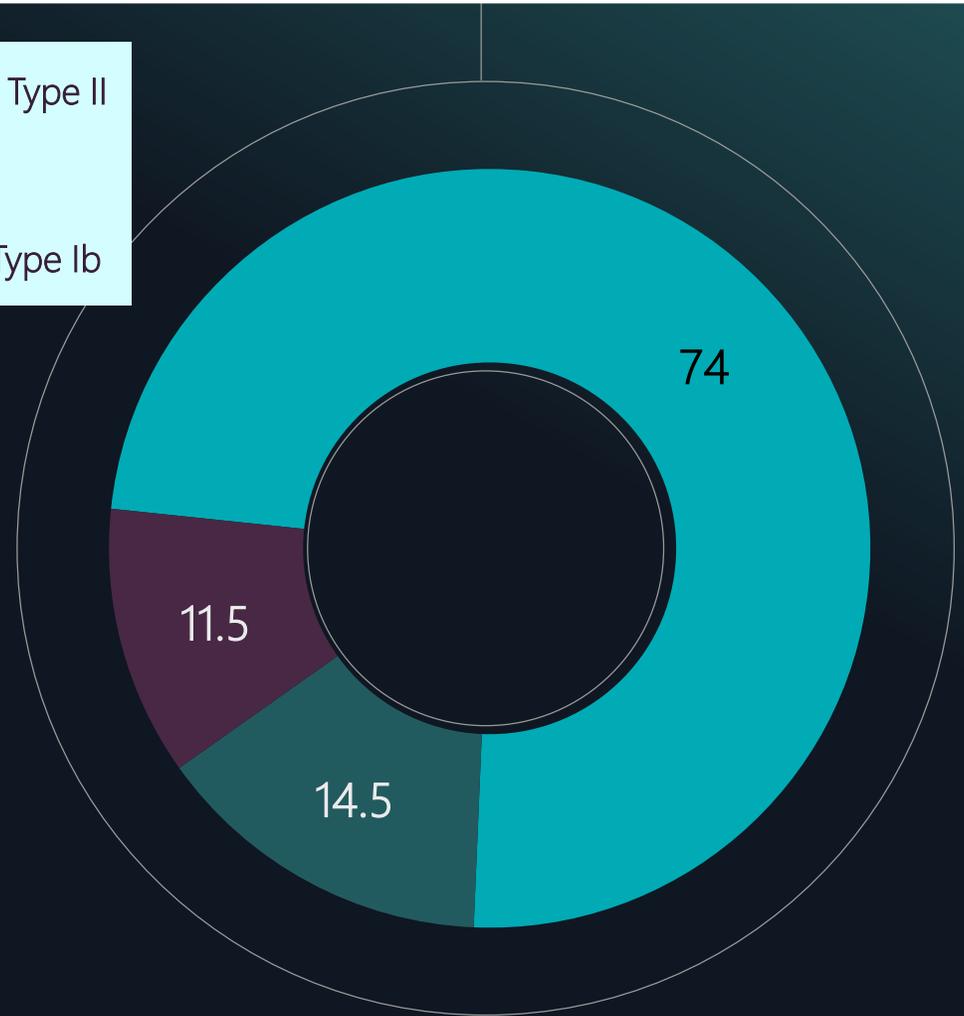
The first step is to select models which

- (a) produce black holes or neutron star remnants, and
- (b) are H and He poor (Ic SN progenitors)

Type II
H & He rich

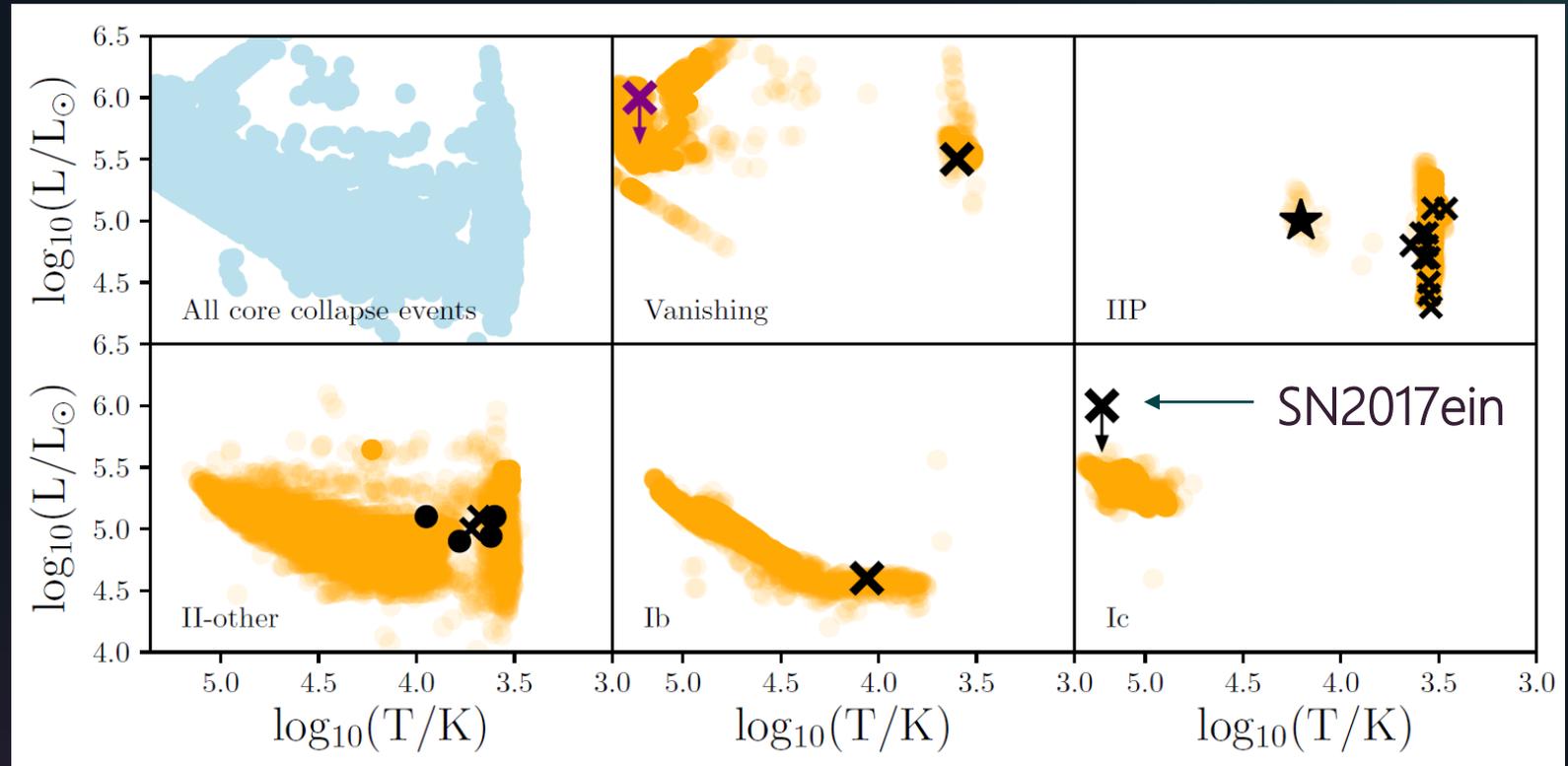
Type Ib
H poor & He rich

Type Ic
H & He poor



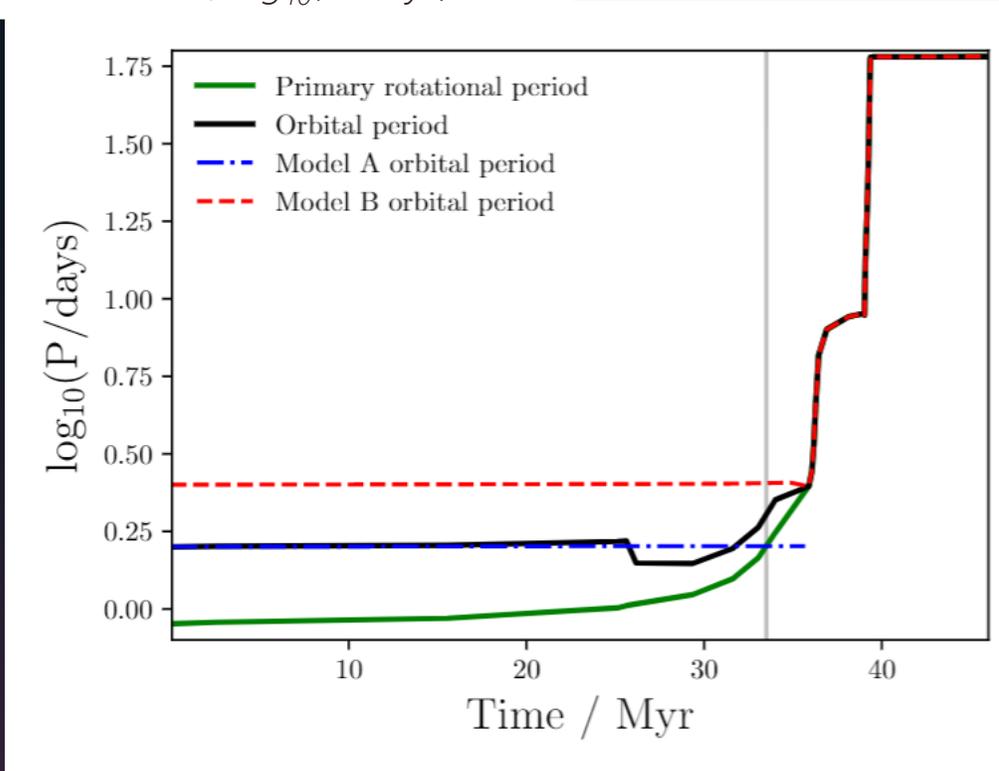
Progenitors on the HRD

- BPASS models
- Pre-explosion progenitor imaging



Tidal post-processing

$8M_{\odot}$, 0.9 mass ratio, $\log_{10}(P/\text{days})=0.2$



We need to consider stellar spins: the models are post-processed to include tidal interactions¹.

Binaries are synchronised upon primary envelope expansion, and kept tidally locked after that.

MCMC

Four parameters are varied to fit the observed GRB rate as a function of redshift.

$$J_{\text{cut}} = J_{\text{cut},\theta} \times \left(\frac{z}{z_{\odot}}\right)^n$$

Choosing GRB models



$J_{\text{cut},\theta}$

The minimum angular momentum required for GRB production.

z^n

The dependence of J_{cut} on metallicity. The index n is varied.

Intrinsic to observed



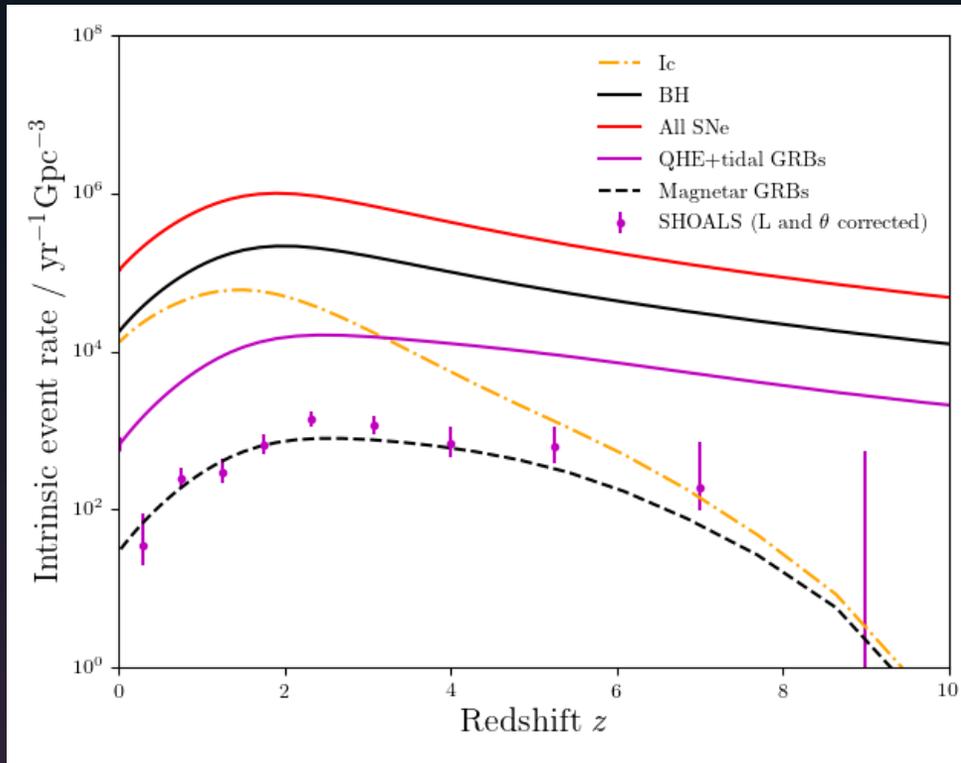
θ

The jet opening angle. Corrects the intrinsic rate by a factor $[1-\cos(\theta)]^{-1}$.

E_{min}

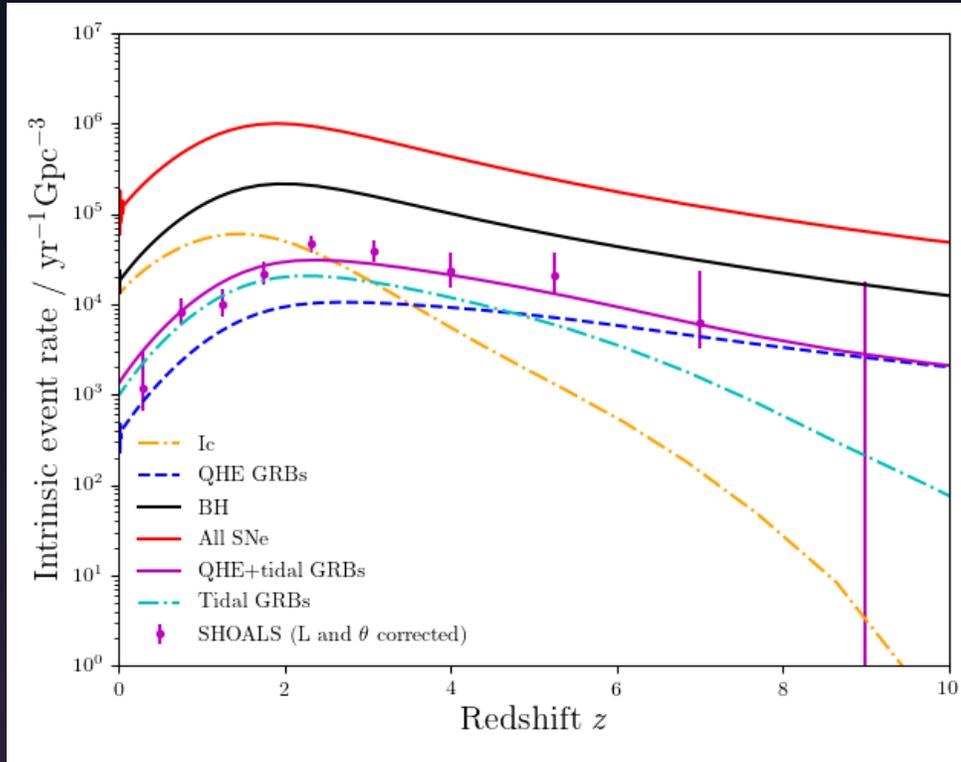
The minimum energy that GRBs release. A luminosity function is assumed with E_{min} as the lower bound.

MCMC Results - Magnetars



	Median
θ /deg	$12.8 \pm_{4.6}^{6.4}$
$\text{Log}_{10}(E_{\text{min}}/\text{erg})$	$49.6 \pm_{0.5}^{0.6}$ 
$\text{Log}_{10}(\dot{J}_{\text{cut},\theta} / \text{cm}^2\text{s}^{-1})$	$18.8 \pm_{0.4}^{0.3}$
Z^n index	$2.2 \pm_{1.3}^{1.0}$

MCMC Results - Collapsars



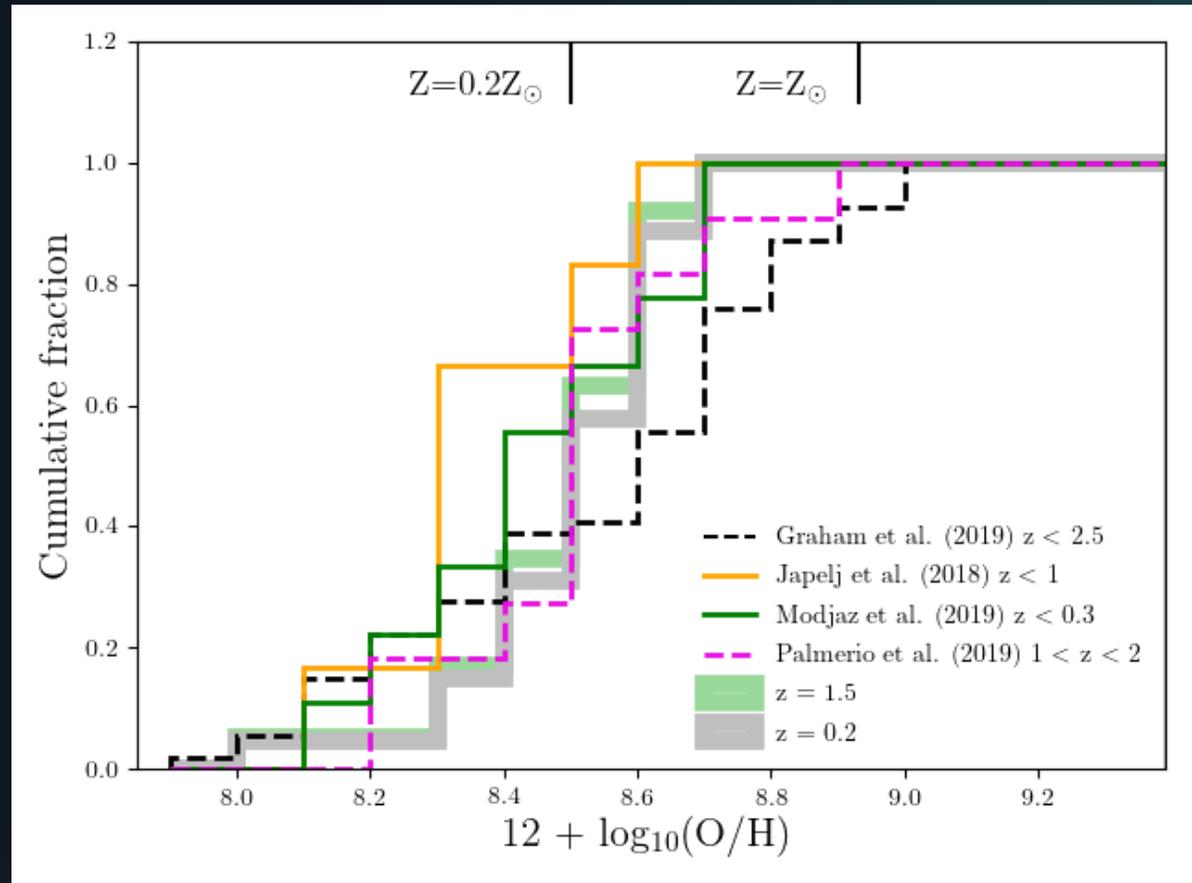
	Median
θ /deg	$9.9 \pm_{6.8}^{8.9}$
$\text{Log}_{10}(E_{\text{min}}/\text{erg})$	$48.1 \pm_{0.7}^{1.2}$
$\text{Log}_{10}(\dot{J}_{\text{cut},\theta} / \text{cm}^2\text{s}^{-1})$	$18.7 \pm_{0.6}^{0.4}$
Z^n index	$3.8 \pm_{2.4}^{2.1}$



Metallicity distribution

The metallicity distribution, assuming collapsars and an $n=3.8$ metallicity dependence.

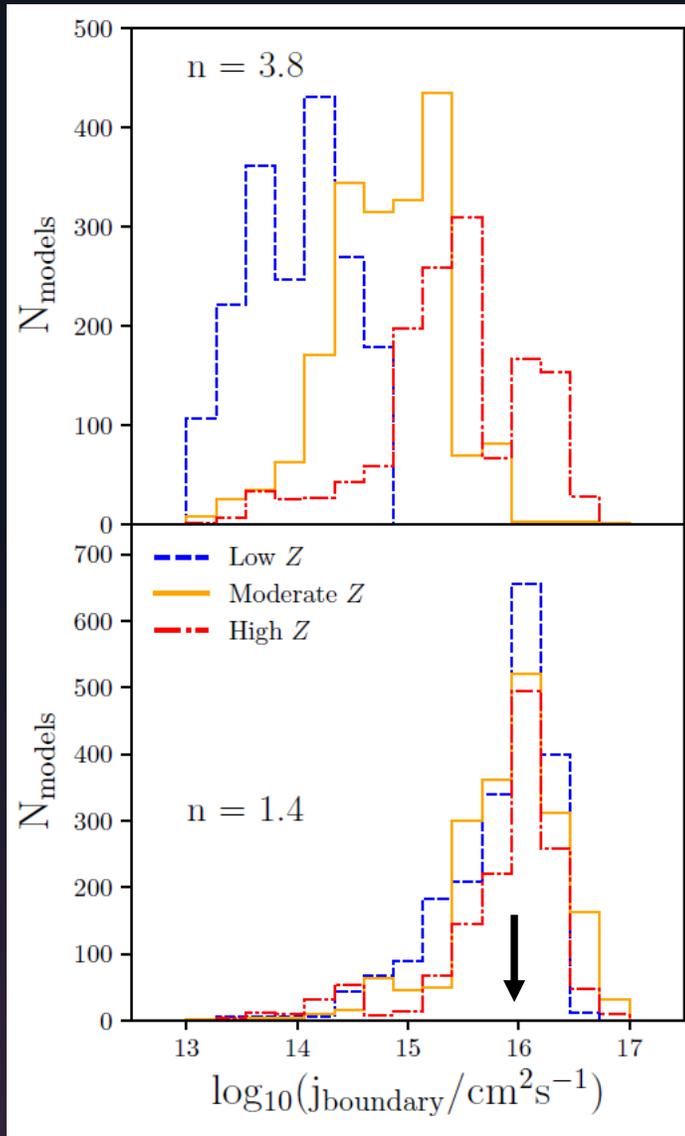
We can also predict the T and L of the progenitor stars prior to core-collapse.



Core angular momentum

Using the STARS model outputs, the accretion disc j can be calculated assuming rigid body rotation, and that the accretion disc forms at the radius enclosing the subsequent remnant mass.

$$j_{\text{cut}} = j_{\text{cut},0} \times \left(\frac{Z}{Z_{\odot}}\right)^n$$



If the 1σ lower bound for n is used, the core j required agrees across the range of metallicities (no reason why that should be Z dependent?)

$\sim 10^{16} \text{ cm}^2\text{s}^{-1}$ is the minimum required in the collapsar model!

$\rightarrow Z^n$ where $n \sim 1$ is also approximately how opacity scales due to metals in a stellar envelope – higher opacity – higher j_{cut} required



SUMMARY

Binary population synthesis with BPASS can reproduce the long GRB rate and host galaxy properties if:

- There are two channels, QHE and tidally spun-up,
- Black holes are the dominant engines,
- There is a metallicity dependence on the escape of relativistic jets from the stellar envelope once they are launched

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Any questions?

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